## Beam Physics Note 6



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# Present situation of the cleaning insertions IR3 and IR7

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The lattices described here are matched to LHC version V6.-2. The main quadrupole strengths used in the next two sections are:

### Main quads:

KQF := .0087696KQD := -.0085463

## 1 IR3

The preferred layout is:

- Q6 cold
- 2 polarities of warm quadrupoles MQW, i.e. addition of a symmetric element to the MQW's
  - standard main quads in the DS

#### Reasons:

- 1) Q6 cold is the only solution which provides DXN=  $D_x/\sqrt{\beta_x}$  above 0.2  $m^{1/2}$  (0.22  $m^{1/2}$  reached, Fig 1). Q6 is composed of a standard cold main quad and 2 trims of length 1.15 m.
- 2) The addition of a symmetric element to the MQW's is mandatory to maintain the strength of QT12.L3 below the limit and the strength of the MQW's below 0.0015  $m^{-2}$ .
- 3) With an exact antisymmetry of the MQW's (without symmetric elements) the trim strengths can not be reduced to acceptable values even by introducing MQL in the IR3 DS.

Fig 1 shows an IR3 lattice with Q6 cold and with symmetric MQW's (quad strengths KQ6S.3, KQ5S.3, KQ4S.3). The corresponding table shows the K1 values of the Q6-trims, warm and DS-trim quads, in this order. There are 24 warm modules, of them 4 of the

F and D remind you the nearest main quad polarity. A higher DXN peak is impossible to get because KQT12.L3 exceeds  $0.0047 \ m^{-2}$  (110 T/m).

Fig 2 shows IR3 lattice with Q6 warm and symmetric MQW's. The location of the DXN (=0.18) peak is tuned to be located between the dogleg magnets.

Fig 3 shows the same case of warm Q6, but without symmetric MQW's (set to zero). Here KQT8.L3, KQT10.L3, KQT12.L3 (all "F") exceed twice the allowed limit (=0.0047). For the first two, increasing the corresponding main quad lengths from 3.1m to 3.45m will not help because this would contribute only 0.00877 \* 0.35/1.15 = 0.0027 to the K1 values. The length of the main quad Q12, cannot be increased.

## 2 IR7

The preferred layout is:

- Q6 cold
- 2 polarities of warm quadrupoles MQW, i.e. addition of a symmetric element to the MQW's
  - standard main quads in the DS

Reasons:

- 1) Q6 cold is the only solution which provides a dispersion D<25 cm in the straight section. With Q6 warm, we get  $D \sim -100$  cm at some location, while the maximum allowed is  $\sim 40$  cm.
- 2) having Q6 cold and addition of a symmetric element to the MQW's provides a larger asymmetry near the middle of the insertion, which is an important ingredient for efficient collimation. Such an increased asymmetry is more important than a net increase of the total available phase advance interval (as in case of Q6 warm).

Simulations using the code DJ show that the collimation quality can be related to the shape of the tune-split function measured from the location of the primary collimators:

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\mu_y(\mu_x) (\mu_y = \mu_x = 0 at the primary), or equally: (\mu_y - \mu_x)/2 depending on (\mu_y + \mu_x)/2.
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The four primary jaw-collimators are located within the drift of the left dogleg.

The value of the escaping amplitudes is related to the shape and the position of the curve  $\mu_y(\mu_x)$  with respect to the optimum contours  $A(\mu_x, \mu_y) = 1$  corresponding to a secondary amplitude normalized to a secondary collimator aperture. In case of Q6 cold and in presence of symmetric modules this shape can be made more irregular by increasing the symmetric parts: KQ4S and KQ5S. Having such an irregular shape and the ability to control it is the preferable option.

In the limiting case of exactly regular  $(\mu_y - \mu_x)/2$ , the phases above 0.5 (×2 $\pi$ ) are not used, because the conditions for collimation repeat itself with the periodicity of the contours

Fig 4 shows the case Q6 cold and symmetric modules. The tune interval available for secondary jaws is  $0.002 < \mu_x < 0.587$ . The results are:

16 jaws: 
$$A_{max} = 8.9$$
,  $A_{x,max} = 8$ ,  $A_{y,max} = 7.7$ 

Fig 5 shows the case Q6 warm and symmetric modules. The tune interval available for secondary jaws is larger:  $0.002 < \mu_x < 0.8$ . The result is:

16 jaws: 
$$A_{max} = 8.9$$
,  $A_{x,max} = 8.8$ ,  $A_{y,max} = 8.4$ 

Fig 6 shows the case Q6 warm and symmetric modules set to zero.

Fig 7 shows the case Q6 cold with symmetric modules set to zero. The result is:

16 jaws: 
$$A_{max} = 9.2 A_{x,max} = 8.35 A_{y,max} = 7.87$$

These surviving amplitudes are larger than in the case of Q6 cold and nonzero symmetric modules.

## 3 Lattices matched to arc with cell tunes

$$\mu_{x,cell} = 0.28 = \frac{7}{25}, \ \mu_{y,cell} = 0.24 = \frac{6}{25}$$

The main quadrupole strengths are:

Main quads:

KQF := .0091824KQD := -.0084275

Two cases were considered:

Fig 8 (/ir3/q6cold) and

Fig 9 (/ir7/q6cold).

DJ analizys is not yet done, but the optics in both cases looks adequate for colliation.

## 4 Conclusions

In both IR3 and IR7, the requested performance of the insertion to do good collimation implies

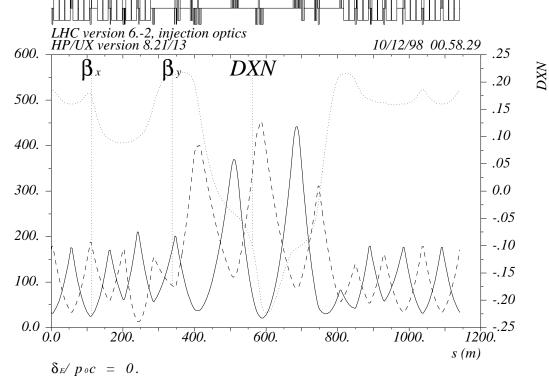
- The use of a cold Q6 made of a regular arc quadrupole (MQ,3.1m long)) and of two TRIM's 1.15m long.
- The use of an additional 'symmetric' MQW (same hardware, electric cabling modified), for a total number of 24 MQW modules in each insert ion.

With this option, regular MQ (3.1 m long) and short trim's (1.15 m long) can be used in the dispersion suppressors.

In IR3, the present limit towards better performance is the strength of MQ12.

In IR7, the present limit towards better performance is the strength of both MQ12 and MQ7, the latter one being easy to overcome (space fo r a second trim exists).

Would Q6 be warm, the number of MQW modules would be 38 in IR3 and 42 in IR7.



 $Table \ name = TWISS$ 

Fig 1: /ir3/q6cold

 $\beta$  (m)

KQT6.R3	.0030790
KQT6.L3	0024657
KQ5A.3	0012874
KQ5S.3	.0011948
KQ4A.3	.0012231
KQ4S.3	.0010519
KQT12.L3	.0043824
KQT11.L3	.0027619
KQT10.L3	0000679
KQT9.L3	0036215
KQT8.L3	0011485
KQT7.L3	0012736
KQT11.R3	0033986
KQT10.R3	.0008716
KQT9.R3	0000325
KQT8.R3	.0033000
KQT7.R3	.0008228

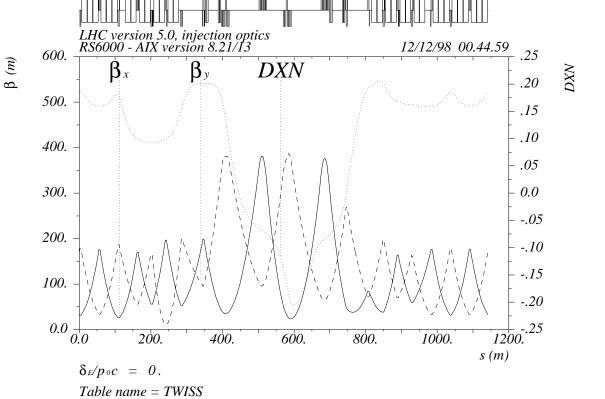


Fig 1a: /ir3/q6cold with aperture limitations: betamax at Q5A.R3 reduced by 40m; DXN = 0.2

KQT6.R3	0.0033112
KQT6.L3	-0.0021544
KQ5A.3	-0.0012681
KQ5S.3	0.0011446
KQ4A.3	0.0012360
KQ4S.3	0.0008785
KQT12.L3	0.0039481
KQT11.L3	0.0024340
KQT10.L3	0.0005570
KQT9.L3	-0.0044559
KQT8.L3	-0.0004571
KQT7.L3	-0.0014935
KQT11.R3	-0.0035092
KQT10.R3	0.0028607
KQT9.R3	-0.0000706
KQT8.R3	0.0035988
KQT7.R3	-0.0004883

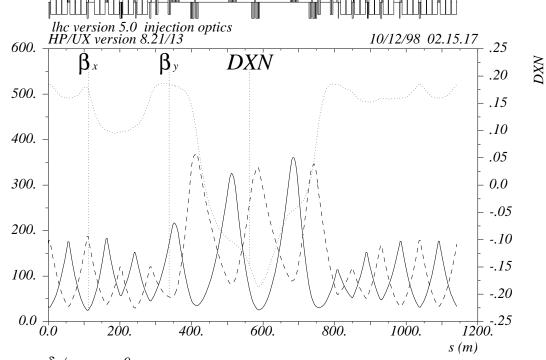
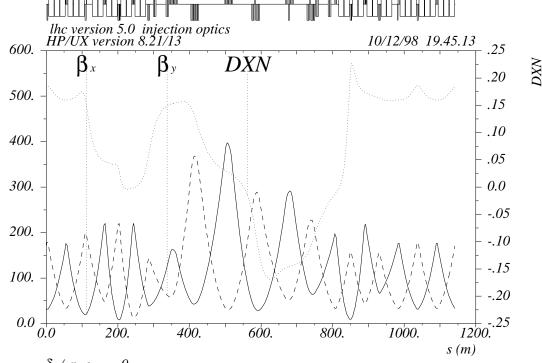


Fig 2: /ir3/q6warm

 $\beta$  (m)

.0015480 0005325	> limit
.0000020	
0014881	near limit
.0007761	
.0012423	
.0010473	
.0004229	
.0008176	
0020552	
0020158	
.0011837	
0022599	
.0004091	
0007749	
.0048854	
0018727	
	.0005325 0014881 .0007761 .0012423 .0010473  .0004229 .000817600205520020158 .0011837 0022599 .00040910007749 .0048854



 $\beta$  (m)

Fig 3: /ir3/q6warm with small or zero symmetric modules (KQ4S.3, KQ5S.3 and KQ6S.3)

KQ6A.3 KQ6S.3 KQ5A.3 KQ5S.3	.0015688 .0000357 0014437 0000051	> limit
KQ4A.3 KQ4S.3	.0013369 .0000443	
rd±b.o	.0000440	
KQT12.L3	.0096773	> limit
KQT11.L3	.0016176	
KQT10.L3	.0118267	> limit
KQT9.L3	0046816	
KQT8.L3	.0111512	> limit
KQT7.L3	0043206	
KQT11.R3	0041398	
KQT10.R3	0008230	
KQT9.R3	.0040913	
KQT8.R3	0030939	
KQT7.R3	.0043593	



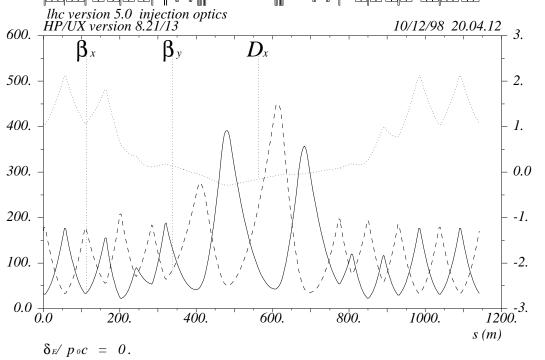


Fig 4: /ir7/q6cold

KQT6.R7 KQT6.L7 KQ5A.7 KQ5S.7 KQ4A.7	.001264600114840012706 .0002222 .0012553
KQ4S.7	.0000416
KQT12.L7	0007046
KQT11.L7	.0025927
KQT10.L7	.0045950
KQT9.L7	.0003270
KQT8.L7	.0014896
KQT7.L7	.0026132
•	
KQT11.R7	.0002057
KQT10.R7	.0004712
KQT9.R7	.0036480
KQT8.R7	.0028023
KQT7.R7	.0031240

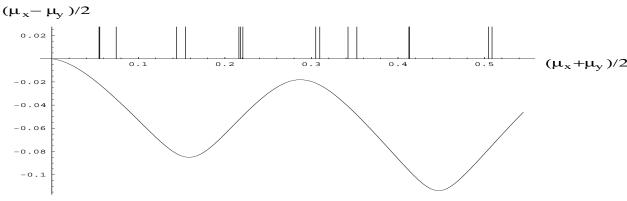


Figure 4 a) The split function  $(\mu_y - \mu_x)/2$  vs  $(\mu_y + \mu_x)/2$  for the lattice on Figure 4 (ir7/q6cold). At the primary collimator we choose to have  $\mu_x = \mu_y = 0$ . The setup of 16 sec. jaws presented on the plot by vertical lines corepond to  $A_{max} = 8.9$ ,  $A_{x,max} = 7.95$ ,  $A_{y,max} = 7.7$ 

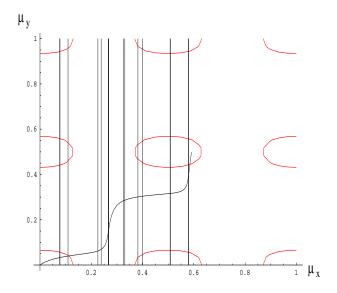


Figure 4 b) The split function, same as on the above plot, but on the plane  $\mu_x, \mu_y$ . At the primary collimator  $\mu_x = \mu_y = 0$ . Collimation quality can be related to the position of the split curve with respect to the contours  $A(\mu_x, \mu_y) = 1$ , also shown.



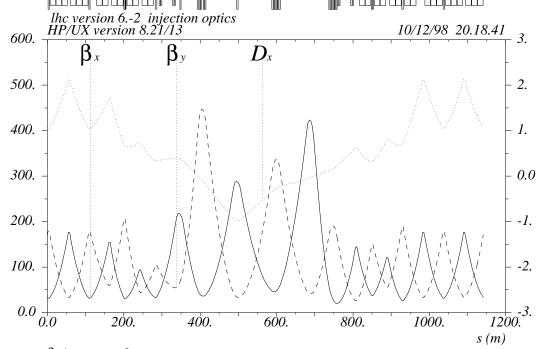


Fig 5: ir7/q6warm

KQ6A.7	.0017914	>	limit
KQ6S.7	0001735		
KQ5A.7	0015492	>	limit
KQ5S.7	.0002803		
KQ4A.7	.0012207		
KQ4S.7	0000509		
KQT12.L7	.0001529		
KQT11.L7	.0028953		
KQT10.L7	.0031035		
KQT9.L7	0009310		
KQT8.L7	.0034114		
KQT7.L7	.0044823		
KQT11.R7	.0007981		
KQT10.R7	.0000633		
KQT9.R7	.0023029		
KQT8.R7	0025714		
KQT7.R7	.0041087		

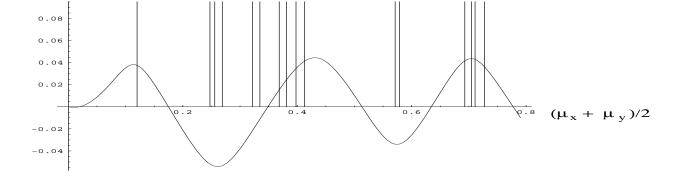


Figure 5 a) The split function  $(\mu_y - \mu_x)/2$  vs  $(\mu_y + \mu_x)/2$  for the lattice on Figure 5 (ir7/q6warm). At the primary collimator we choose to have  $\mu_x = \mu_y = 0$ . The setup of 16 sec. jaws presented on the plot by vertical lines corespond to  $A_{max} = 8.84$ ,  $A_{x,max} = 8.8$ ,  $A_{y,max} = 8.4$ 

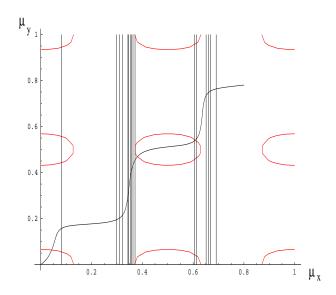


Figure 5 b) The split function, same as on the above plot, but on the plane  $\mu_x, \mu_y$ . At the primary collimator  $\mu_x = \mu_y = 0$ . Collimation quality can be related to the position of the split curve with respect to the contours  $A(\mu_x, \mu_y) = 1$ , also shown.



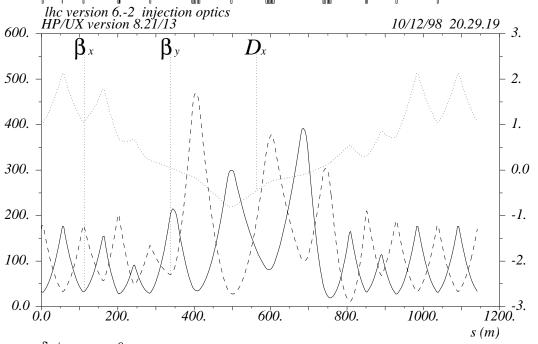


Fig 6: /ir7/q6warm with small or zero symmetric modules (KQ4S.7, KQ5S.7 and KQ6S.7)

KQ6A.7	.0018006	> limit
KQ6S.7	.0000088	
KQ5A.7	0014955	near limit
KQ5S.7	.0001063	
KQ4A.7	.0010854	
KQ4S.7	.0000000	
KQT12.L7	0004300	
KQT11.L7	.0026983	
KQT10.L7	.0033843	
KQT9.L7	0006624	
KQT8.L7	.0053666	
KQT7.L7	.0060610	> limit
KQT11.R7	.0005081	
KQT10.R7	.0015058	
KQT9.R7	.0027048	
KQT8.R7	0067280	> limit
KQT7.R7	.0050084	near limit



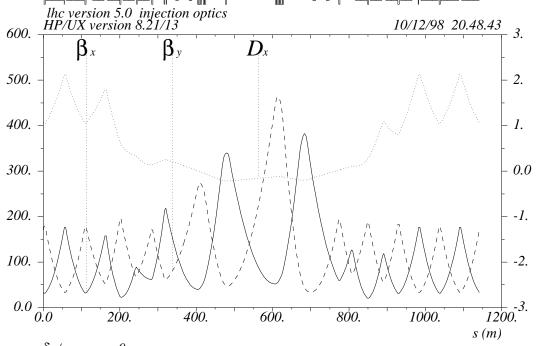
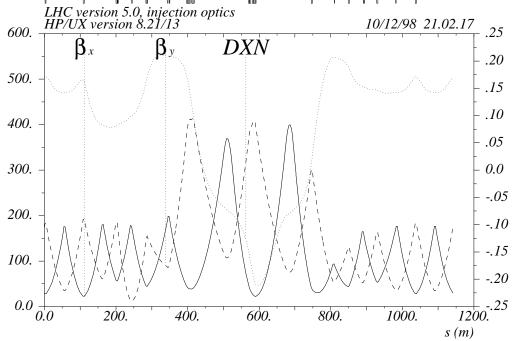


Fig 7: /ir7/q6cold with small or zero symmetric modules (KQ4S.7 and KQ5S.7)

KQT6.R7 KQT6.L7 KQ5A.7 KQ5S.7 KQ4A.7 KQ4S.7	.001001700130080012691 .0000323 .00126000000027
KQT12.L7	0003196
KQT11.L7	.0022548
KQT10.L7	.0046966 near limit
KQT9.L7	.0003270
KQT8.L7	0000308
KQT7.L7	.0025880
KQT11.R7	0000297
KQT10.R7	.0002257
KQT9.R7	.0038695
KQT8.R7	.0026197
KQT7.R7	.0032651





DXN

 $\delta E/p_0c = 0.$ 

Table name = TWISS

Fig 8: /ir3/q6cold; arc cell tunes mux=0.28, muy=0.24

.0026032
0025578
0013058
.0011737
.0012187
.0009719
.0042789
.0026561
.0001054
0038791
0016799
0020801
0034836
.0005650
0007408
.0028728
0008644



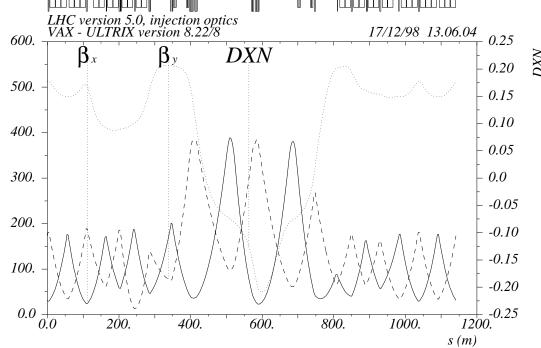
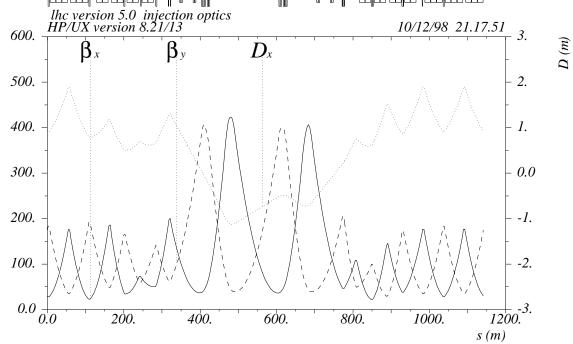


Fig 8a: /ir3/q6cold; Talman tunes: Qx=65.28, Qy=58.31

KQT6.R3	0.0029018
KQT6.L3	-0.0022577
KQ5A.3	-0.0013134
KQ5S.3	0.0010343
KQ4A.3	0.0012440
KQ4S.3	0.0008523
KQT12.L3	0.0037336
KQT11.L3	0.0028937
KQT10.L3	-0.0000870
KQT9.L3	-0.0034769
KQT8.L3	-0.0012166
KQT7.L3	-0.0025476
KQT11.R3	-0.0036473
KQT10.R3	0.0023001
KQT9.R3	-0.0005135
KQT8.R3	0.0033794
KQT7.R3	-0.0009632





 $\delta_{E}/p_{0}c = 0.$ 

 $Table\ name = TWISS$ 

Fig 9: /ir7/q6cold; AV tunes: arc cell tunes mux=0.28, muy=0.24

KQT6.R7 -.0002568 KQT6.L7 -.0011952 KQ5A.7 -.0013527 KQ5S.7 .0001526 KQ4A.7.0013158 KQ4S.7 .0000275 KQT12.L7 .0041601 KQT11.L7 .0014327 KQT10.L7 .0031769 KQT9.L7 .0000904 KQT8.L7 -.0039819 KQT7.L7 .0035451 KQT11.R7 -.0015264 KQT10.R7 -.0034732 KQT9.R7 .0024965 KQT8.R7 -.0011394 KQT7.R7 .0036008