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Collimation efficiency in the presence of collimator misalignment and sample closed orbit errors

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Summary

We present robustness studies of the LHC collimation system and report tracking results at injection obtained with a combined code DIMAD + STRUCT. Besides a description of the simulation tool, featuring inelastic nuclear scattering and chromatic tracking including sextupoles, we discuss first results on combined betatron and momentum collimation inefficiency when the off-momentum halo is included. We assume (as a simplified example) a constant aperture limit $R = 2$ cm around the ring at all drift entrances, a sample closed orbit error generated by transverse offsets of two arc quadrupoles and corrected within the insertions, and random transverse collimator misalignments.

1 Introduction

We have used Dimad ([2]), to track for several hundreds of turns the halo created at injection by the IR7 primary vertical collimator. The LHC lattice is version V6.3 for Beam 1, with thick elements. Tracking includes particle propagation through the collimator material, a feature made possible by the new collimator element in Dimad, based on the main block of the code STRUCT, [1]. This tool should be considered complementary to other available tools for collimation studies (K2, SIXTRACK). The new capabilities of Dimad are described in the Appendix.

At present, a matter of great concern is the behaviour of an imperfect collimation system, and in particular how to set tolerances on jaw misalignment and orbit excursions within the collimation insertions. A model of a collimation system with errors has therefore been developed and first results from tracking are presented here, and also compared (section 4.2) with previous results from other approaches [4]. Some of the model features and simplified assumptions are: it uses the full lattice at injection (both IR7 and IR3 insertions); it assumes a constant around the ring radial aperture; it applies sample closed orbits outside the collimator sections and random transverse collimator misalignments within; it takes into account the energy losses occurring as halo particles traverse the jaws. To explain our choice of model and the approximations made, the following considerations apply.

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At injection, one is mostly interested in halo loss rates in the arcs. The inefficiency is therefore computed with respect to arc losses by setting an aperture limitation $(R=2 \text{ cm})$ at all drift entrances. By shifting two arc sextupoles, an artificial error closed orbit is created. We correct the orbit to zero within the collimation insertions, but impose large collimator misalignment with respect to the chamber axis (this should give a rough idea of what happens if the orbit is not corrected).

No field errors have been included, the chromaticity being set to 2 units in both planes by the lattice sextupoles alone. The code Dimad is based on second order TRANSPORT maps with kicks describing the action of higher order multipoles and also accepts symplectic ray tracing [3]. It is believed that it describes with sufficient accuracy the short-term motion $(< 1000 \text{ turns})$, including off-momentum and far from the axis.

The jaw material and thickness used are: for primary $-$ Al (20 cm); for secondary $-$ Cu (50 cm).

The off-momentum halo created by energy losses in the jaws

At injection, the initially nearly on-momentum particles emerging from the IR7 primary collimator lose their energy via inelastic nuclear interactions within the jaws with resultant negative momentum offsets of the circulating halo down to about -1% .

The motion within the collimator material is governed by the STRUCT block, an important parameter being the maximum loss of momentum that a particle may suffer during nuclear interaction. By setting the corresponding STRUCT parameter dpplim to (practically) infinity, we assume that all off-momentum particles are kept.

We will show that under such assumption: 1) the losses on the 2-cm radial aperture are largely caused by off-momentum tertiary halo particles; 2) in the case of a realistic and hence misaligned system IR7 and IR3 work together even with respect to halo born at the IR7 primary collimator only.

Secondary halo consists of particles emerging from the primary collimator during the first passage. These do not reach the arc or the IR3 collimators in one turn, but are either absorbed during the first passage through IR7, or receive additional deflecting kicks on subsequent passages and are gradually absorbed or scattered in the IR7 secondary collimators,

Due to the nature of nuclear interactions, the particles emerging at some turn from the IR7 secondary collimators (tertiary halo) may acquire both large negative momentum offsets and large betatronic amplitude. Such particles are lost at the high dispersion locations in the arc, or reach the high-dispersion place in IR3 and hence the IR3 collimators (mostly during the same turn).

At the same time, the capacity of IR3 for stopping off-momentum particles effectively is restricted by the value of the ring momentum acceptance $|\delta_{cut}| = 3.5 \; 10^{-3}$. This is because IR3 is designed [∗] to effectively absorb the **circulating beam** particles, i.e. the ones with $|\delta| < |\delta_{cut}|$. The above criterion however can only serve as a rough estimate.

When designing the ideal system (for $\Delta p/p = 0$), it has been assumed that the total halo created at the IR7 primary collimators is almost entirely absorbed in the IR7 secondaries. This indeed turns out to be the case, with the exception of the highly off-momentum tertiary halo described above.

[∗]By increasing the normalized dispersion at the IR3 primary collimator.

2 Initial conditions and tracking procedure

Optics, LHC lattice, collimator settings. This study is based on injection optics Version 6.3 (tunes $Q_x=64.28$, $Q_y=59.31$) with the emittance and beam energy being set to either their injection or collision values. The IR7 and IR3 collimators are adjusted to their ideal apertures, according to the corresponding beam envelopes (see below the definition of ideal system).

The choice of the **equilibrium impact distribution** (see [7], [8]) is crucial for defining efficiency.

The initial point for tracking is set just before the first primary (vertical) collimator of IR7[∗] and two beams, each composed of particles with zero normalised angle and zero momentum offset, are generated entering the upper and lower jaw (Fig. 1). Each beam has zero size in direction parallel to the jaw surface and a rectangular distribution over impact parameters. We use an elimination procedure – only particles impacting within 1 micrometer from the jaw edge are tracked. The corresponding BEAM and GENERATE commands are shown in the Appendix. The two beams are

Figure 1: (colour) The first (vertical) primary collimator and the impact distribution (red) $-1 \mu m$ size in y direction and zero size in direction parallel to the jaw border.

tracked until all particles are lost in collimators, or on the aperture limitations. At injection, even for the ideal system, less than 200 turns are needed. Tracking is stopped, if less than 10 particles are still circulating. Whether this small but long-living halo fraction will reach the aperture, or will be captured, does not change the results.

3 Collimation systems with misalignment

We consider collimators fixed at some, possibly imperfect, positions with respect to the vacuum chamber axis. In a machine with misalignments, the on-momentum closed orbit within the IR7 and IR3 insertions deviates from this axis (kind 1, Fig 2). Among many possible corrected systems, such as: fully corrected within the insertions (kind 3), orbit corrected to some extent at the primary collimator only $(kind 2)$ – both may be with or without collimator misalignment – we only consider the following case:

orbit corrected within the collimator-occupied sections (same as if the collimators were aligned w.r.t. uncorrected orbit) plus random transverse displacement of the midpoint (center) of each pair jaws w.r.t. to the vacuum chamber axis. All collimator centres are randomly displaced, except for the first primary collimator in IR7.

Table 1 defines the four different cases considered:

[∗] A similar procedure and lattice has been set to model the halo created by IR3 (not reported here)

orbits

Figure 2: Possible uncorrected and corrected Figure 3: The orbit $CO2 - 4 \text{mm}$ peak to peak in the arc

	Table 1: Collimation systems with misalignment
ideal	all collimators at ideal apertures no misalignment
corrected CO ₂	orbit corrected to zero within the COS sections
	corrected $CO2 + jaw$ mis. corrected plus transverse collimator misalignment

- ideal system is the error-free machine plus sextupoles (at chromaticity 2), with all collimators set at their ideal normalized aperture w.r.t. the vacuum chamber axis: $n_1 = 6, n_2 = 7$ in IR7, $n_1 = 8.5, n_2 = 9.5$ in IR3 (σ units; n_1 for primary, n_2 for secondary) and with rotation angles as defined in the LHC database [∗] . The collimator parameters of the ideal system, namely lattice functions at collimators and jaw orientation angles, are the same as in www.cern.ch/LHC-collimation.
- By shifting a pair (QF and QD) of arc quadrupoles in transverse directions, a nonzero closed orbit is created, with nearly equal r.m.s. and maximum displacements in both planes (Table 2 and Figure 3). For the **corrected** system the orbit is corrected to zero within the collimatoroccupied sections, i.e. from the entrance of the first to the exit of the last collimator in each IR7 and IR3. In practice, such correction is achieved by applying four combined (coordinate and angle) kicks (see Fig 4). The kicks preserve the uncorrected orbit outside the two COS. We notice that 1) the orbit within the IR7 Right Dispersion Suppressor (RDS7) is thereby left uncorrected; 2) an ideal every-turn correction has been assumed.
- For the corrected $+$ jaw mis. the midpoints of the collimators are assigned random transverse displacements dx, dy with respect to the vacuum chamber axis. A flat distribution is used within the interval $(-|dx|_{max}, |dx|_{max})$ and the same for y.

3.1 The "2-cm inefficiency"

By setting a radial aperture R=2 cm at all drift entrances, the "2-cm inefficiency" w.r.t. particles lost in some section of the ring is defined to be the ratio:

 N_{part} lost at 2cm-apert. section/Ntotal lost part.

[∗] The real-space distance between the jaw and the chamber axis is expressed through the emittance, horizontal and vertical beta function at the collimator, its rotation angle and the normalized aperture.

	CO ₂
MAX at arc quadrupoles	
RMS at arc quadrupoles	1.4
MAX at IR7 collimators	1.5

Table 2: Uncorrected $CO2$ – approximate excursions in mm (nearly equal in x and y)

Figure 4: The two lattice sections occupied by collimators in IR7 and IR3 called below collimator-occupied sections COS7 and COS3. Four orbit-correcting kicks are applied at the entrances and exits of the two COS. The region outside both COS is left uncorrected – it consists of RDS7, ring 7-3 and ring 3-7 (RDS7 is excluded from ring 7-3).

The four different aperture sections of interest are COS7, COS3, outside both COS and ring 7-3 + ring 3-7. The last (fourth) inefficiency part corresponds to particles lost outside both COS and also escaping the RDS7 (we denote it by " ring 7337").

4 Results

4.1 The effect of the variable dpplim

If during a nuclear interaction the momentum offset of a particle $|\Delta p/p|$ exceeds dpplim, the particle is counted to be lost in the collimator. Particles reaching high betatron amplitudes also have large momentum offset $|\Delta p/p|$, hence the choice dpplim strongly affects the fraction reaching the 2-cm aperture.

For all calculations dpplim is set to the maximum value 0.3, as allowed by STRUCT. In effect, all off-momentum particles are allowed to circulate freely (without acceleration). With dpplim set to 0.5%, all losses are strongly reduced, as shown in Table 3 for one seed.

	lnı		l ot	
DPPLIM=	0.3	.005	0.3	.005
absorbed in coll-s	0.9447	0.9676	0.9988	0.9994
\vert COS7 aperture	$0.0538 \mid 0.0314$			
\vert COS3 aperture		0.0004		
outside COS aperture	$0.0013 \mid 0.0009$		0.0008	

Table 3: Effect of changing dpplim. Fraction absorbed in collimators and losses on the 2-cm pipe. Numbers smaller that 5.10[−]⁵ are shown as zero. The IR3 collimators are retracted.

4.2 Integrated inefficiency and comparison with the results in [4]

By removing the 2 cm aperture limitation, the integrated inefficiency found by DIMAD + STRUCT has been compared with the one in [4] (collision energy and optics; IR7 collimators only; micrometer impact depth on the first primary, uniform distribution of impact parameters). The corresponding inefficiency curves are shown on Figure 5 – the lower curves, corresponding to 1μ m initial beam size are to be compared with the results in [4]. The seed to seed variation is seen in Fig. 4.2.

Sample one-seed runs, with and without aperture limitation, were made at injection and collision. Table 4 shows the integrated inefficiency value at 20 sigma (collision and injection) and the 2-cm inefficiency computed over different ring sections.

	integr. ineff. for $A_r = 20$		2-cm inefficiency				
no apert. besides coll-rs		(apert. limit 2 cm))					
	at $D_x \approx 0$	at $D_{x,arc}$	$\overline{\text{in} \text{ coll-s}}$ COS7 COS3 outside COS ring 7337				
Collision							
Ideal	510^{-4}	0.001	0.9989			0.0007	$\sim 10^{-5}$
Injection							
Ideal	510^{-4}		0.9461	0.0523		$.0015\,$	

Table 4: Ideal system (no IR3 collimators) for one seed (10^5 particles) . Left: integrated inefficiency as defined in [4]; Right: "2-cm inefficiency" w.r.t. different ring sections.

Figure 5: Ideal system (no IR3 collimators) for one seed $(8\ 10^4\ \text{particles})$. Integrated inefficiency at collision (as in [4]). Left: computed at location just before the first primary; Right: at maximum dispersion near one of the arc focusing quadrupoles.

Figure 6: Seed-to-seed variation of the integrated inefficiency shown in Figure 5, left.

4.3 Effects of orbit error and jaw misalignment at injection

The fraction of halo reaching the ring 7337 aperture is of the order of 10^{-4} , hence for meaningful statistics one needs to track around 10^6 particles. We choose instead to track several smaller patches, each corresponding to a new initial seed for the impact distribution and simultaneously new initial seed in the STRUCT module, and record average and maximum values over the patches. With $8\,10^4$ particles per patch, the patch to patch fluctuation is due solely to the STRUCT random generator, since the impact distribution is sampled with sufficient accuracy. For the present study, ten such patches (seeds) have been generated and applied to each Ideal and Corrected CO2 system.

The 2-cm aperture limitation was applied for all calculations in the section. Initially, the IR3 collimators were retracted. The first two rows of Table 5 show average and maximum values of the quantity in each column with an accuracy down to 10 particles ($\sim 10^{-5}$). For the average (first row), the sum of the numbers in columns 2 to 5 is unity. The fraction lost in RDS7 is given by the difference of the numbers in the last two columns.

As seen in Table 5, the presence of the corrected CO2 causes an increase of the average losses over ring 7337 by 35 %, w.r.t Ideal, while the maximum value changes half as much.

	Halo losses, injection, no IR3 coll-s					
	10 seeds $\times 810^4$ part					
	absorbed in	on the 2 cm radial apert				
	collimators	COS7		$COS3$ outside COS ring 7337		
	Ideal					
ave	0.94514		0.05344 0.00002	0.00140	1.710^{-4}	
max	0.94606	0.05424	0.00005	0.00156	2.510^{-4}	
	$Ideal + corrected CO2$					
ave	0.94493	0.05355	0.00002	0.00149	2.310^{-}	
max	0.94604	0.05407	0.00003	0.00160	2.810^{-4}	

Table 5: Average and maximum values of 10 seeds (STRUCT random generator). The IR3 collimators are retracted

Halo losses, injection, no IR3 coll-s						
worst seed w.r.t. CO2 effect						
	absorbed in			on the 2 cm radial apert		
	collimators	outside COS ring 7337 COS7 COS3				
Ideal	0.94474	0.05383	θ .	0.00135	1.610^{-4}	
corr. CO ₂	0.94494	0.05355	θ .	0.00155	2.810^{-4}	
corr. $CO2 +$ worst	0.94389	0.05412	0.00005	0.00205	7.610^{-4}	
seed 0.5 mm jaw mis.						
corr. $CO2 +$ worst	0.94385	0.05158	0.0001	0.00455	1.910^{-3}	
seed 1 mm jaw mis.						

Table 6: The worst seed w.r.t. CO2 effect on ring 7337 losses.The IR3 collimators are retracted.

Some explanations apply. For all ten random systems, when the corrected orbit was added, the fraction lost outside the collimator occupied section increased, while the fraction lost in ring 7337 was sometimes smaller or remained the same – the increase being absorbed inside RDS7 (where the orbit is not corrected). Another factor is the action of the orbit correcting kicks, which tend to improve quality. They redirect the particles at the entrance to the primary collimator, thus sending them into the secondary jaws a little more efficiently.

Introducing the IR3 collimators was found to have an improving statistical effect, with respect to the maximum inefficiency values, i.e. on the worst seeds. Table 6 shows the worst seed with respect to effect of the corrected CO2, i.e. the seed that produced the largest increase of ring 7337 occupancy after corrected CO2 was added (this seed turned out to be also the worst w.r.t. increase of losses outside COS). For this seed, the jaw misalignment was added on top of corrected CO2 (Table 6) and the IR3 collimators were switched on – Table 7. Their improving effect w.r.t ring 7337 loses is seen to be around 60%.

With jaw misalignments, the total fraction escaping COS and reaching the rest of the ring is determined by the partially destroyed collimation quality (larger escaping window through the collimators), as shown in Table 8. For jaw misalignment amplitudes 0.5 mm and 1 mm, the ring 7337 losses increase 4.7 and 12 times respectively. In this case the IR3 collimators have been unable to improve efficiency, as is seen in Table 7.

Halo losses, injection, with IR3 coll-s					
worst seed w.r.t. CO2 effect					
on the 2 cm radial apert					
ring 7337 outside COS					
Ideal	0.00137	1.410^{-4}			
corr. CO2	0.00150	1.710^{-4}			
corr. $CO2 +$ worst	0.00205	7.410^{-4}			
seed 0.5 mm jaw mis.					
corr. $CO2 +$ worst	0.00443	1.810^{-3}			
seed 1 mm jaw mis.					

Table 7: Same seed as on Table 6, but IR3 collimators included.

Individual particles are lost on the COS3 aperture even if the IR3 collimators are missing – this is because of the high-dispersion peak still being present there. If IR3 was missing together with its optics, these would have been lost in the arcs.

The fraction lost over ring 7337 is not only composed of particles reaching the arcs, but also includes the losses at the high beta location in IR6, element BPMYC (where radius 2 cm corresponds to aperture $\sim 10 \sigma$). The fraction lost on BPMYC was between 30 and 70 % of the total ring 7337 occupancy and for most seeds increased when errors were included.

Table 8: Halo losses for the 10 seeds of $CO2 + jaw$ misalignment $(|dx|_{max} = |dy|_{max} =$ 1 mm).

The code provides a listing of the names of all elements receiving non-zero losses and the number of particles lost. The following listing corresponds to the ring 7337 case of Table 6 with 0.5 mm jaw misalignment.

5 Summary and conclusions

Comparisons of the integrated inefficiency, as computed with the DIMAD + STRUCT model, with previous results [4] (collision energy, no aperture limitations) show satisfactory agreement [6].

We have studied the inefficiency of an imperfect collimation system at injection. Only the halo produced at the IR7 primary collimator has been simulated. We have approximately replaced the LHC vacuum chamber with a 2-cm radial aperture limitation at all drift entrances.

For the ideal system, studied with $\sim 10^6$ particles, the 2-cm inefficiency with respect to ring 7337 was found to average 1.7 10^{-4} with maximum 2.5 10^{-4} (this includes the losses on a ~ 10 sigma absorber) and w.r.t. to the RDS7 – $\sim 1.310^{-3}$ ^{*}.

A sample closed orbit, corrected at every turn to zero within the jaw sections of both IR7 and IR3, has been found to increase the average occupancy of ring 7337 by around 40 %, nearly preserving the maximum occupancy. Introducing IR3 has in this case an improving effect on the worst error seeds (around 60 %).

The worst seed of random uniform transverse misalignment of all collimators in both planes and both directions with amplitude 0.5 mm, added on top of the corrected orbit, increased the losses nearly 5 times with respect to Ideal.

In future one should consider a more realistic chamber geometry. Correct treatment of the off-momentum particles requires acceleration to be included in the model.

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References

- [1] I. Baichev, D. Kaltchev, Implementation of a new Collimator Element (STRUCT module) in Dimad for the purposes of LHC Collimation, TRIUMF Report (in preparation)
- [2] R.V. Servranckx, K.L. Brown, L. Schachinger, and D. Douglas, User's Guide to the Program DIMAD, SLAC Report 285 UC-28 (A), May 1985 (http://www.triumf.ca/people/kaltchev/dimad.html)
- [3] D. Douglas, Etienne Forest, R.V. Servranckx, A Method to Render Second Order Beam Optics Programs Symplectic, IEEE Transactions of Nuclear Science, Vol. NS-32, No 5, October 1985.
- [4] R. Assmann, J.B. Jeanneret, D. Kaltchev, Status of the Robustness Studies for the LHC Collimation, Proc. APAC01 (http://www.cern.ch/lhc-collimation)
- [5] R. Assmann et al, "Preliminary Beam-based Specifications for the LHC Collimators", LHC Project Note 277, 2002 (http://www.cern.ch/lhc-collimation)
- [6] (to be published)
- [7] Thys Risselada, The Effect of the Betatron Tune on the Impact Parameter at the LHC primary collimators, SL Note 92-16 (AP), 1992 (http://www.triumf.ca/people/kaltchev/shelf)

[∗] . The losses occurring in the dispersion suppressor have been studied in more detail in [9].

- [8] M. Seidel, Dissertation, The Proton Collimation System of HERA, DESY 94-103, June 1994 (http://www.triumf.ca/people/kaltchev/shelf)
- [9] I. Baichev, J.B. Jeanneret, Beam losses in the dispersion suppressors of IR3 and IR7, LHC-PROJECT-NOTE-253, 2001.
- [10] U. Wienands, C.P. Parfitt, F.W. Jones, Collimator Simulation for the Triumf KAON Factory using Dimad, TRI-PP-91-42, May 1991 and Particle Accelerator Conf., San Francisco, CA, May 6-9, IEEE 1991.

6 Appendix: collimation capabilities of Dimad

To study propagation of the halo, the basic block of code STRUCT has been implemented in Dimad. Next we describe the new collimator element and set-collimator operation, discuss some universal and special (LHC) features, and present sample input and output lists.

Tracking of GeV protons with their penetration and scattering inside the collimator material was first introduced in Dimad during the KAON factory studies [10]. The collimator is treated during tracking as an arbitrary element [2]. In a similar way, for the purposes of LHC collimation, a new beamline element (rectangular collimator) has been created that coincides with the basic module of STRUCT. The element has three parameters: aperture [m], rotation angle [rad] and kind of material (3 possible kinds). The parameters are passed to the tracking module via a new operation SETCOL, which can be used along with the standard Dimad operations such as tracking, field errors and misalignment, closed orbit correction, analysis of geometric and chromatic aberrations etc.

The input lattice for Dimad is, with minor differences (such as as TITLE), identical to the one of MAD8 (line format). Just as in MAD, the USE command is applied to identify the beamline. Next follow DIMAT, to pass the expanded line to the Dimad routines, the BEAM, to define the initial beam centroids (an arbitrary number of them) and corresponding beam sizes, and GENERATE – to generate the desired particle number and distribution. At this stage commands related to errors and misalignment can be introduced.

Before one proceeds with tracking, the SETCOL command must be called at least once for every collimator present in the beamline of use, in order to indicate how individual collimators are to be treated. The parameters of SETCOL allow:

– when applied to a primary collimator, to trigger the "discard by impact" procedure, i.e. to discard at the first turn coordinates getting outside some desired impact distance. By adjusting the parameters of BEAM and GENERATE, an arbitrary impact distribution may be defined.

– to treat the collimator as a drift, or as a black absorber;

– to store in disk file coordinates or/and amplitudes of particles reaching the collimator entrance;

In accordance with the standard database names of primary and secondary collimators of IR3 and IR7, the kind of collimator is recognised by the combination of characters ("P" or "S") and (3 or 7) in the collimator name.

By installing an arbitrary number of test collimators with arbitrary names and applying to them appropriate SETCOL operations, the halo can be analysed at chosen locations around the ring.

To this end, the beam line must contain a special collimator element named DUMMY, preceding any other collimator and invisible to the beam (with large aperture and zero length). This element serves to pass, via corresponding SETCOL DUMMY operation, global parameters of the collimation system and *global commands*, i.e. such that will affect all collimators during one or more tracking runs. Global parameters are:

– the beam energy and emittance;

– the ideal normalised apertures for primary and secondary collimators in IR3 and IR7 (given these, all collimators are intrinsically adjusted to their real-space offsets and angles according to the emittance);

– the momentum cut-off DPPLIM;

The global commands switch on or off:

- the radial aperture at the entrance of all drifts;
- the random transverse misalignment;

– a procedure that distributes particles in amplitude or coordinate bins, as they pass an arc quadrupole (not applied in this paper).

Sample Dimad input:

```
! Snapshot date: 03/12/01 Time: 131444
... <-- LHC B1 lattice here in line format
                                           STARTS AT IR7 PRIMARY called TCPA
k2d = -0.106404818276 <-- must have correct chromaticity 2
k2f=0.065210125333
TCPA: COLLIMATOR,L=.05,p1=1,p2=0.006,p3=1.57
... p1 = material
... p2 = half dist between jaws
                                            p3 = angle in rad
USE, LHCB1<br>MATRIX
                                        \leftarrow show the 2 ord matr. and chrom.
MACHINE
                                        \leftarrow- the Twiss
...
SET COL
DUMMY
SET COL<br>TCPA 5 0 10 10 0,
                         \leftarrow-- name, iprint, iblack, Apx, Apy, m/turn
...
SET SYMPLECTIC OPTION ON BEAM
                             BEAM <-- beam sigmas in TRANSPORT not.
1.e-40 0 0 0 0 0
1.e-40 0 0 0 0
1.e-6 0 0 0 <-- mum impact distr (uniform)
1.e-40 0 0
1.e-40 0
1.e-40
-1,
MISALIGNMENT DEFINITION <-- shift in transv. dir. 2 arc quads RQFA78Bx .75e-3 0. 0 0. 0. 0. 0. 0. 1
RQFA78Bx .75e-3 0. 0 0. 0. 0. 0. 1
RQDA78Bx 0 0. .75e-3 0. 0. 0. 0. 1 CO2 = 4 mm peak to peak
99,
SET MISALIGNMENT<br>REFERENCE ORBIT
                                          REFERENCE ORBIT <-- find closed orbit
...
GENERATE PARTICLES
45000 <-- per jaw
0.0 0.0 0.66224207E-02 0.69429323E-04 0.0 0.0 <-- beam centroids
0.0 0.0 -0.66224207E-02 -0.69429323E-04 0.0 0.0,
TRACKING
```
The Dimad printout contains:

- total number of particles before and after discard by impact;
- collimator names and number of particles lost in each collimator total (third column) and within turn intervals: $0 - 10$, $10 - 20$, ... (next five columns);
- numbers of particles and fractions lost in collimators and on the aperture inside and outside the collimator occupied sections.

Sample Dimad output

```
OPERATION LIST ,
TRACKING
tot part 90000
total latt elts 11517
tot arc quads 392
   turn ncpart
   1 90000
discarded (impact) = 15279
    2 9437
...
   65 6
LAST 6 ARE NOT TRACKED
Loss statistics after passage 65
    coll. name lost 10 20 30 40 50
   1 DUM 0 0 0 0 0 0
   2 TCPA 29044 28640 214 171 14 5 <-- first primary
   3 TCPB 8 8 0 0 0 0 of IR7
   4 TCPC 718 679 31 6 1 1
   5 TCPD 831 789 22 17 2 1
   6 CS71 4974 4909 37 26 1 1
  7 CS72 3086 3033 30 18 4 1
   8 CS73 3569 3519 24 23 3 0
   9 CS74 6083 5994 47 34 6 2
  10 CS75 6336 6253 48 31 2 2
  11 CS76 4822 4750 36 29 5 2
  12 CS77 1382 1363 12 7 0 0
  13 CS78 1710 1680 17 11 2 0 <-- all IR7 secondary<br>14 CS79 52 52 0 0 0 0 0 coll-s are loaded
            52 52 0 0 0 0 0 coll-s are loaded
  15 CS710 3251 3204 27 17 3 0
  16 CS711 319 315 1 3 0 0
  17 CS712 43 43 0 0 0 0
  18 CS713 295 273 18 4 0 0
  19 CS714 180 170 7 3 0 0
  20 CS715 3585 3522 38 22 2 1
  21 CS716 329 324 2 2 0 1
  22 TCP3 0 0 0 0 0 0
  23 CS31 0 0 0 0 0 0 0 <-- the IR3 coll-s. are retracted 24 CS32 0 0 0 0 0 0
  24 CS32 0 0 0 0 0 0
  25 CS33 0 0 0 0 0 0
  26 CS34 0 0 0 0 0 0
  27 CS35 0 0 0 0 0 0
  28 CS36 0 0 0 0 0 0
  29 TJ 0 0 0 0 0 0
Discarded at impact 15283
Total entered 74717
Collimator losses
IABS=1 25593
Win 44778
DppLim 246
total lost in coll-s 70617 0.94513
```
